



# Low-loss GeO<sub>2</sub> thin films deposited by ion-assisted alternating current reactive sputtering for waveguide applications



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## ABSTRACT

We report low optical loss germanium oxide (GeO<sub>2</sub>) thin films which have been deposited by low pressure ion-assisted alternating current dual magnetron sputtering of germanium targets in an oxygen plasma environment. The germanium oxide films ranging from 0.7- to 1.0- $\mu\text{m}$ -thick were fabricated at low temperature and high deposition rates of 6–38 nm/min on silicon and thermally oxidized silicon substrates. The refractive index of the films was determined through variable angle spectroscopic ellipsometry and found to be 1.6051 on average at a wavelength of 638 nm. The films were shown to be near stoichiometric through Rutherford backscattering spectroscopy analysis and displayed sub-nanometer roughness when measured with atomic force microscopy. High peak-to-valley uniformity on a 3-inch substrate with relative deviation of 0.8% was achieved. We characterized the attenuation of the GeO<sub>2</sub> thin films from visible to near-infrared wavelengths and observed it to be as low as 0.1 dB/cm at 638 nm for deposition rates of 8 nm/min. This deposition technique provides complementary metal-oxide-semiconductor compatible GeO<sub>2</sub> thin films suitable for integrated photonics applications and is promising for the fabrication of other dielectric waveguide materials.

## 1. Introduction

Germanium oxide (GeO<sub>2</sub>) is a material that is widely used in photonic applications because of its advantageous structural and optical properties. GeO<sub>2</sub> first gained interest in its applications in doped optical fibers due to its structural similarities to silica, allowing for ease of integration into fiber cores [1]. GeO<sub>2</sub> glass has strong inter-ionic forces between the Ge<sup>4+</sup> and O<sup>2-</sup> ions which provide high thermal stability and mechanical strength [2]. It has a higher refractive index ( $n = 1.6$ ) than silica and is used to increase the numerical aperture of optical fibers. GeO<sub>2</sub> possesses a high Raman gain coefficient when compared to other oxide-based glasses [3], making it a strong candidate for Raman based amplification. It also shows promise as a host for rare earth dopants amplifiers and lasers due to its high rare-earth solubility [4]. Its low nonlinearity and absorption allow for long-distance optical transmission in fibers [5]. GeO<sub>2</sub> glass possesses strong opto-refractive effects allowing for permanent shifts in refractive index when exposed to UV radiation [4–7]. This property is used to write Bragg gratings in optical fibers by UV laser exposure for applications, including fiber lasers, sensors, and dispersion compensators [8–11].

GeO<sub>2</sub>'s high refractive index and low loss also make it a promising material for planar waveguide applications. Because it is a

complementary metal-oxide-semiconductor compatible material already extensively used in silicon foundries for high-k gate dielectrics [12,13], it is particularly attractive for silicon photonics platforms [14]. Its high rare-earth (RE) solubility motivates its use in on-chip amplifiers and lasers. Optical gain and lasing have been demonstrated in germanosilicate glass [15] and several authors have reported efficient photoluminescence in RE-doped germanate alloys such as lead-, antimony-, and bismuth-germinate glasses [16–20]. The UV sensitivity of GeO<sub>2</sub> suggests it might be useful for grating structures or as a UV-sensitive cladding material for post-process trimming to increase yield in silicon photonics devices [21].

While the physical and material properties of GeO<sub>2</sub> thin films have been explored, there has been limited investigation towards the fabrication of low loss thin films. GeO<sub>2</sub> thin films have been previously fabricated by radio frequency (RF) sputtering and reactive direct current (DC) sputtering [2, 22–24]. Both of these techniques have been able to achieve high deposition rates; however, they have not been able to produce films with as-deposited losses of less than 1.9 dB/cm at 638 nm. The losses in these films have been reduced by post-deposition annealing or laser annealing, which adds processing steps and increases fabrication complexity [2]. Germanate based alloy waveguide structures have achieved losses ranging from around 0.6 to 2.2 dB/cm at

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632 nm, with these losses being attributed to scattering off of the waveguide sidewalls [25,26].

Alternating current (AC) dual magnetron sputtering in conjunction with high energy ionic bombardment is known to yield high-quality thin films with high density and uniformity [27]. While other films such as magnesium oxides and tantalum oxynitrides have been fabricated and reported on for optical applications [28,29], GeO<sub>2</sub> thin film fabrication has not been extensively explored using this deposition technique.

In this paper, we study GeO<sub>2</sub> thin film deposition using ion assisted AC dual magnetron sputtering. We deposit the films in varying regions of the sputtering hysteresis bias curve and study their structural and optical properties. The film stoichiometry is confirmed through high energy ion-matter interactions using Rutherford backscattering spectrometry (RBS) and elastic recoil detection analysis (ERDA). We measure low propagation losses using the prism coupling technique and observe low surface roughness using atomic force microscopy (AFM). These properties, combined with high deposition rates and insensitivity to process variation, demonstrate ion-assisted AC dual magnetron sputtering to be a highly promising method for planar waveguide fabrication.

## 2. Experimental methods

### 2.1. Ion-assisted AC dual magnetron sputtering of GeO<sub>2</sub> thin films

We deposited the GeO<sub>2</sub> thin films using an Intlvac Nanocrhome IV Plasma Assisted Reactive Magnetron Sputtering (PARMS) system. A schematic of the system is displayed in Fig. 1. The tool is a top-down sputter system that can be equipped with up to four 6" sputter sources for pulsed DC or dual AC sputtering. In addition, the system contains an ion source for in-situ pre-cleaning and ion-assisted deposition. The PARMS system is designed for the production of dense and low defect films. Prior to the deposition, the chamber was pumped down to  $1.3 \times 10^{-7}$  kPa with a cryopump. Subsequently, the films were deposited under low heating, at 120 °C, using two 6-inch 99.995% pure germanium targets. The plasma processing parameters were varied across oxygen flow ranging from 40 to 80 standard cubic centimeters per minute (sccm) at constant power. In order to maintain film

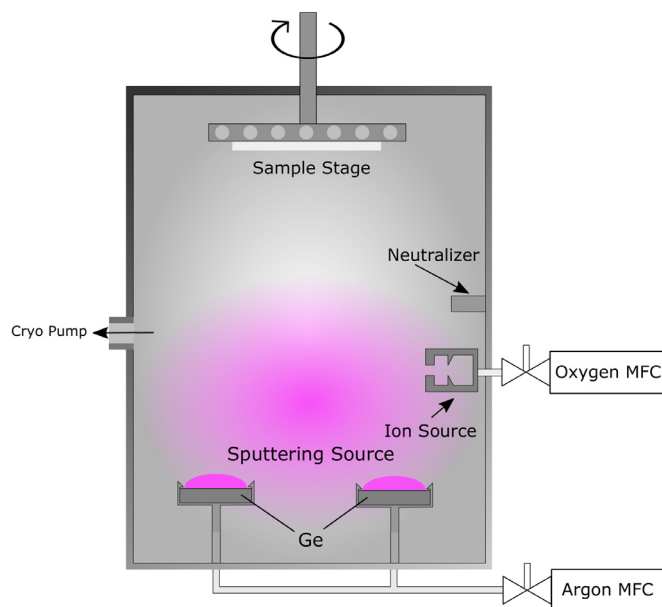


Fig. 1. Schematic of the ion-assisted dual AC sputtering system applied to deposit GeO<sub>2</sub> thin films. The position of the ion source and neutralizer are not representative of the actual position and the system is not drawn to scale.

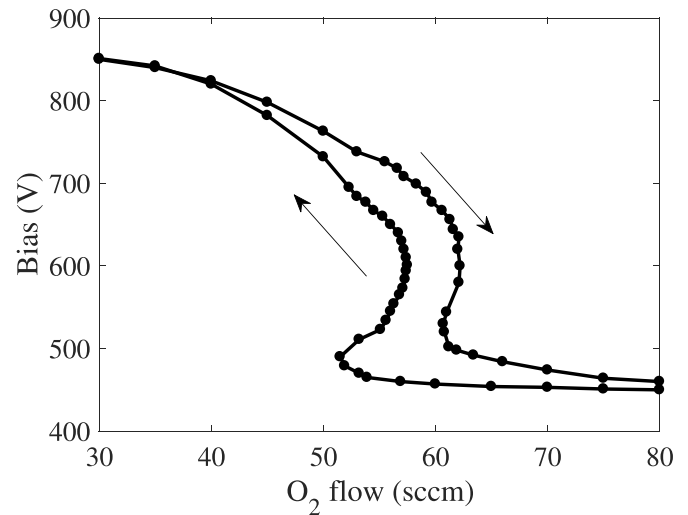


Fig. 2. Target bias hysteresis for varying oxygen flow under ion-assisted deposition.

homogeneity during the growth process, the oxygen flow was either held constant or adjusted by an algorithm to maintain a selected target bias. In order to select the deposition parameters across different sputtering regimes, the target bias response for the system was examined. The target bias vs. O<sub>2</sub> flow rate hysteresis curves are shown in Fig. 2. The transition from metal to poisoned target mostly takes place between 40 and 65 sccm of O<sub>2</sub> flow. The difference between the curve when transitioning from oxide to metal and vice-versa are likely the result of slow variation in the chamber wall and target surface condition [30]. Films were deposited on either side of the hysteresis, i.e. where the target bias is relatively insensitive to oxygen flow and the process is inherently stable, as well as inside the hysteresis, e.g. at 550 V where the algorithm is necessary to stabilize it [31]. The films were deposited on <100> silicon wafers for compositional analysis, as well as silicon substrates with 6- $\mu$ m-thick thermal oxide for waveguide characterization, all fixed to a rotating sample holder.

### 2.2. Thin film characterization

We determined the composition of the GeO<sub>2</sub> films using RBS and ERDA. The RBS results were collected using a 1.8 MeV <sup>4</sup>He<sup>+</sup> ion beam in a Cornell geometry with a dose of 2  $\mu$ C. The sample underwent random rotation during data acquisition to limit the effects of channeling in the crystalline silicon substrate. The ERD results were collected with a 2.9 MeV <sup>4</sup>He<sup>+</sup> ion beam with an incident and exit angle of 75° to the normal of the sample surface. Modeling of the experimental data was performed using the SIMNRA software package to determine the stoichiometric composition of the thin films [32].

We determined the thicknesses and refractive indices of the thin films with the use of a J. A. Woollam M-2000 UV-Vis variable angle spectroscopic ellipsometer (VASE) fitted with a near-infrared extension. Measurements in the spectral range of 250 to 1700 nm were taken at angles of 55°, 60°, 65°, 70°, and 75° from the surface normal. The spectra were then analyzed with the use of J. A. Woollam's CompleteEASE™ software package, assuming a homogenous film comprised of a single layer with no surface roughness on a silicon substrate [33].

In order to determine the optical loss of the deposited material, a set of 0.7- $\mu$ m-thick films were grown onto 6- $\mu$ m-thick thermal oxide layers on 3" diameter Si wafers. The propagation loss of the films deposited on oxidized silicon wafers was measured using a Metricon 2010/M prism coupler at 638, 847 and 1550 nm wavelength. Using a set of linear polarizers, the light was coupled into the film under transverse electric (TE) and transverse magnetic (TM) mode conditions. The mode order was selected by altering the coupling angle of the incident laser onto

the prism with the fundamental TE and TM modes being chosen. Measurements were performed across different regions of the wafer approximately ten times and the statistical average of these measurements was selected as the loss of the film. We carried out multiple measurements to account for the sensitivity of the prism coupling system to the scattering of the propagating light from surface defects along the beam path. Pieces of the GeO<sub>2</sub>-coated oxide oxidized Si wafers were cleaved and annealed separately in an annealing furnace under a constant N<sub>2</sub> gas flow of 60 sccm. The samples were annealed at temperatures ranging from 200 to 400 °C in steps of 50 °C.

The surface roughness of 1.0-μm-thick films on polished silicon was measured using AC mode AFM. A probe stiffness of 60 N/m and a frequency of 300 kHz was used. Three random 2 μm x 2 μm locations on the surface were measured to determine the average root mean squared (RMS) roughness of the film. This scan size was significantly larger than the surface features, thus was determined to be sufficiently large to obtain an average roughness. The data analysis for the AFM measurements was performed in the WSXM 5.0 application [34].

### 3. Results and discussion

Fig. 3(a) shows the experimental RBS spectra and corresponding simulated fit for a GeO<sub>2</sub> thin film grown with an O<sub>2</sub> flow rate of 60 sccm. A single layer model for the deposited film consisting of germanium, oxygen, and silicon was applied to the experimental data. There is no evidence of interstitial argon in the films based on a lack of signal around 1000 keV. From the fitted models, it can be observed that the films are near stoichiometric GeO<sub>2</sub> within the experimental uncertainty for germanium of 0.1 atomic% and oxygen of 0.8 atomic%. This uncertainty is obtained through the fit statistics given by the SIMNRA software. The lower uncertainty in the atomic percentage of germanium compared with oxygen can be attributed to germanium's higher atomic mass, which results in a larger atomic cross-section and a higher probability of scattering incoming He<sup>+</sup> ions. This stronger scattering renders the concentration of heavier elements more sensitive to measurement. The composition of the GeO<sub>2</sub> thin films does not change for different oxygen flow rates operating in the poisoning regime for the Ge

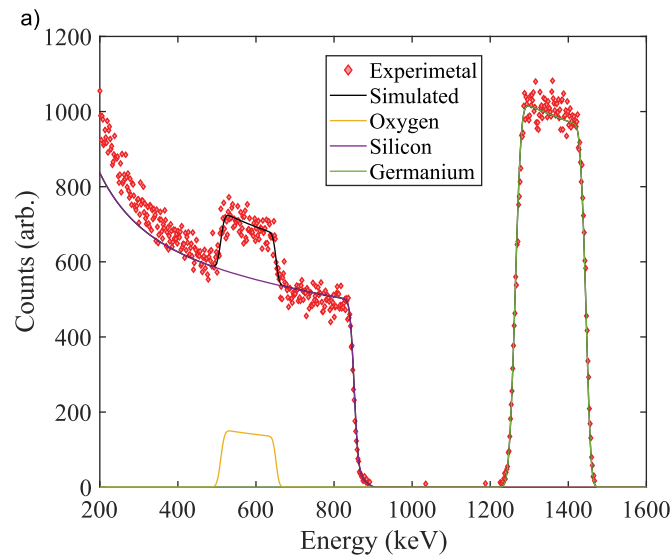


Fig. 3. a) RBS experimental spectra (red dots) and fits based on a simulation model for a GeO<sub>2</sub> thin film. The yellow, purple, and green fit lines represent oxygen, silicon, and germanium, respectively. The composition of the GeO<sub>2</sub> is stoichiometric within the uncertainty in the simulated model. b) ERDA experimental spectra (red dots) with fit based on simulation model for 1.1 at.% hydrogen content uniformly distributed throughout the film. (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article.)

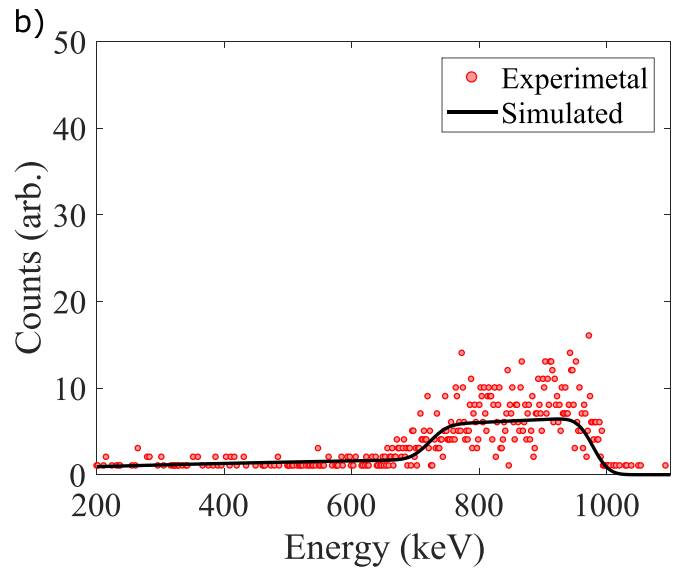


Fig. 3. (continued)

targets as well as the metal sputtering regime and the hysteresis midpoint.

Fig. 3(b) shows the ERDA spectrum measured for the same GeO<sub>2</sub> film in order to characterize the hydrogen content. The simulated model suggests a hydrogen concentration of 1.1 at.% spread evenly throughout the thickness of the film. The source of this hydrogen is not certain. A possible explanation is the instruction of hydrogen into the chamber from residual water vapor which is then incorporated into the film during the deposition. Alternatively, the hydrogen in the film might arise from water vapor diffusing into the film between deposition and ERD measurements. We note that the presence of hydrogen in the films can be useful for specific applications. It is proposed that the presence of OH groups can contribute to defect sites such as germanium electron trap centers and these centers contribute to the photorefractive properties of GeO<sub>2</sub> when exposed to UV light [35,36].

The Cauchy and Sellmeier model were applied to characterize the dispersion properties of the GeO<sub>2</sub> thin films from data obtained through VASE. The film loss was characterized using the Urbach absorption model. For this model, the parameters of the silicon substrate were taken from Herzinger [37]. Fig. 4(a) shows the Cauchy dispersion model and Urbach absorption as well as the Sellmeier model for a 60-

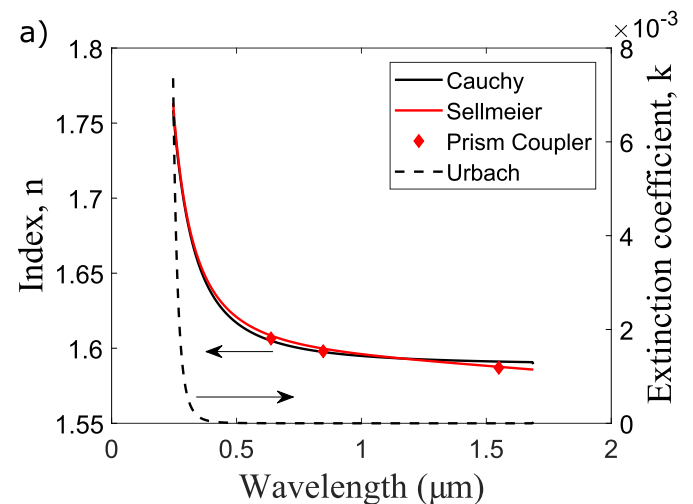


Fig. 4. a) Material dispersion of GeO<sub>2</sub> thin film deposited with 60-sccm O<sub>2</sub> flow measured with VASE and prism coupler. b) Height of GeO<sub>2</sub> thin film in nm using the Cauchy model measured across 3 inch silicon wafer.

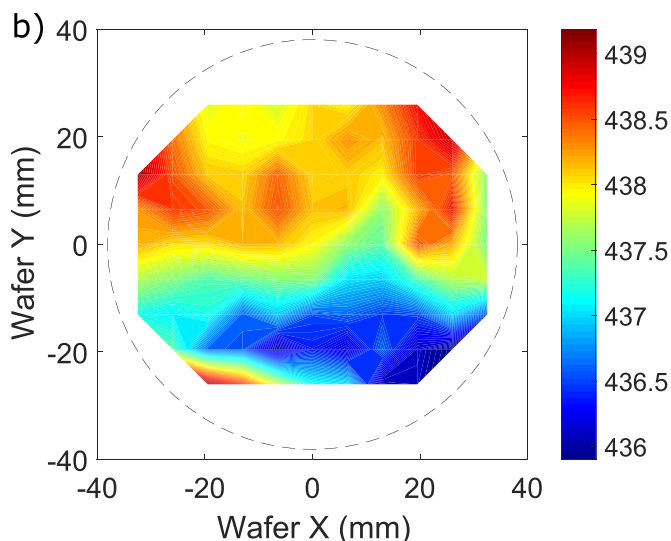


Fig. 4. (continued)

sccm O<sub>2</sub> flow GeO<sub>2</sub> film. Overlaid on top of the dispersion curves are direct refractive index measurements obtained using prism coupling at three separate wavelengths. These overlap with the Cauchy and Sellmeier dispersion model except for the measurement taken at 1550 nm in which the Cauchy model overestimates the index. This difference is expected as the Sellmeier model is considered better over large wavelength ranges. The Cauchy model serves as a reasonable approximation across the visible region and is defined by the three-term function:

$$n = A + B\lambda^{-2} + C\lambda^{-4} \quad (1)$$

where the fitting parameters *A*, *B* and *C* are the material-dependent Cauchy coefficients, *n* is the refractive index of the film and  $\lambda$  is the wavelength. Likewise, a 2-term Sellmeier model was used given by the formula:

$$n^2 = \epsilon(\infty) + \frac{B_1\lambda^2}{\lambda^2 - C_1^2} - E\lambda^2 \quad (2)$$

where  $\epsilon(\infty)$  is the index offset, *B*<sub>1</sub> and *C*<sub>1</sub> are the experimentally determined Sellmeier coefficients and *E* is the position of the pole in the infrared. The values for the dispersion models shown in Fig. 4 are presented in Table 1. The refractive indices reported here are similar to those reported elsewhere using other deposition techniques [2,23]. Fig. 4b) shows the height map of a 3-inch silicon wafer coated in GeO<sub>2</sub> positioned mid-way between the center and the edge of the sample stage. The film shows high uniformity across a large area with a relative deviation in thickness (Max-Min)/average of 0.8% across the 437.7 nm thick GeO<sub>2</sub> thin film with a standard deviation from the mean of 0.2%.

Sputtering at lower oxygen flow rates resulted in a significant increase in the deposition rate, from 8 to 10 nm/min at flow rates of greater than 60 sccm to 38 nm/min at 40 sccm, as shown in Fig. 5(a). This increase in deposition rate primarily can be attributed to the lower O<sub>2</sub> partial pressure, reducing target surface oxidation, thus increasing the germanium sputtering probability at lower O<sub>2</sub> flow rates [31]. There a noticeable trend in the thin film losses deposited at different

Table 1

Cauchy and Sellmeier dispersion parameters coefficients obtained through VASE.

Cauchy	A	B (μm) <sup>-2</sup>	C (μm) <sup>-4</sup>		
	1.587	6.310 × 10 <sup>-3</sup>	2.278 × 10 <sup>-4</sup>		
Sellmeier	ε(∞)	B <sub>1</sub> (μm) <sup>2</sup>	C <sub>1</sub> (μm) <sup>2</sup>	E (μm) <sup>-2</sup>	
	1.665	8.63 × 10 <sup>-1</sup>	1.5447 × 10 <sup>-1</sup>	8.84 × 10 <sup>-3</sup>	

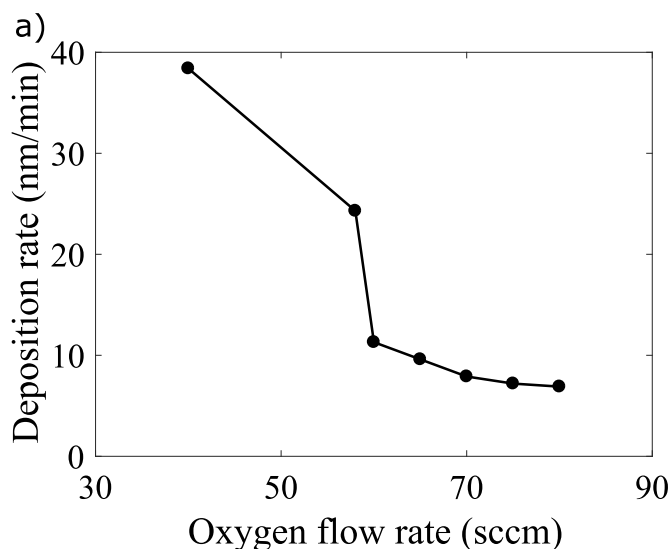


Fig. 5. a) GeO<sub>2</sub> thin film deposition rate vs. O<sub>2</sub> flow rate under ion-assisted growth. b) As deposited film loss at 638 nm measured by the prism coupling method of TE mode light vs. deposition rate. Insert shows 638 nm red streak from prism coupler coupling into GeO<sub>2</sub> film.

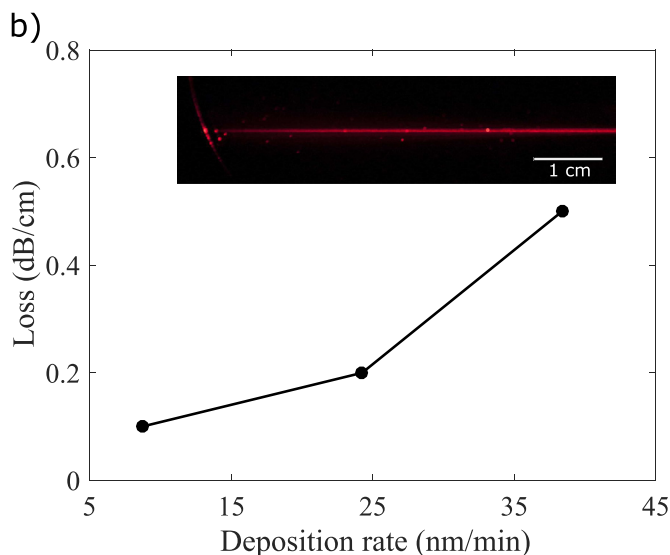


Fig. 5. (continued)

deposition rates. Those films deposited at lower rates in the poisoned regime showed low losses that were at the measurement limit of 0.1 dB/cm at 638 nm across the wafer for both the fundamental TE and TM modes. Films that were grown in the knee of the hysteresis with deposition rates around 25 nm/min showed losses around 0.2 dB/cm at 638 nm, and films grown in the metal sputtering region showed losses around 0.5 dB/cm at 638 nm, for both the fundamental TE and TM modes.

AFM was used to study whether surface roughness could account for the higher losses at 638 nm seen in the films with higher deposition rates. The RMS surface roughness of films deposited with between 60 and 80 sccm of O<sub>2</sub> and with deposition rates below 10 nm/min was typically 0.5 to 0.8 nm. This is consistent with other films grown through PARMs [38]. Annealing selected films in the range of 200 to 400 °C, as shown in Fig. 6, did not yield a noticeable difference in roughness compared to the as-deposited films. It can be noted that there appears to be a correlation between surface roughness and loss, shown in Table 2. The film with an RMS surface roughness of 0.92 nm exhibited the highest loss with 0.5 dB/cm at 638 nm. Conversely, the film

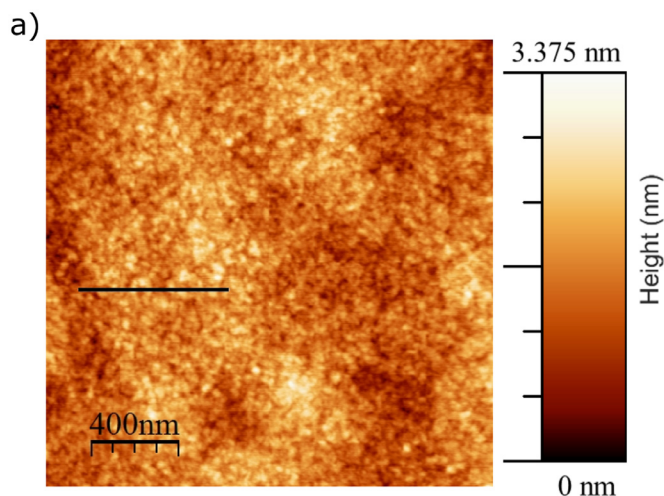


Fig. 6. a) AFM profile for a 1- $\mu\text{m}$ -thick  $\text{GeO}_2$  thin film annealed at 400 °C with RMS roughness of 0.45 nm and peak to peak roughness of 3.8 nm. b) Height profile taken from AFM scan in a).

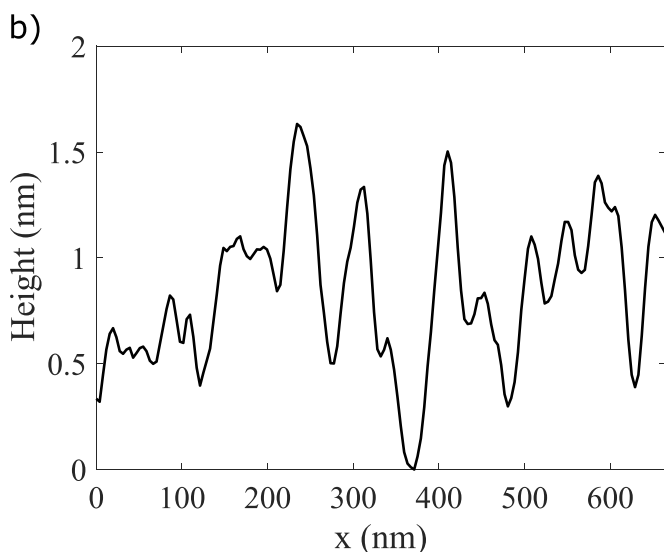


Fig. 6. (continued)

Table 2

Deposition parameters, atomic concentrations obtained through RBS, refractive index obtained through VASE, RMS roughness and peak to peak roughness for films obtained through AFM, losses obtained from prism coupling.

O <sub>2</sub> flow rate (sccm)	Ge (at. % $\pm$ 0.1)	O (at. % $\pm$ 0.8)	N @ 638 nm	Loss, 638 nm (dB/cm)	R <sub>RMS</sub> (nm)	Peak to Peak (nm)
40	31.0	69.0	1.603	0.5	0.92	9.2
55	31.0	69.0	1.600	0.2	0.73	6.2
60	32.5	67.5	1.605	0.1	0.54	4.1
65	32.5	67.5	1.605	-	0.53	5.1
70	32.8	67.2	1.605	-	1.21	8.6
75	33.0	67.0	1.605	-	0.77	8.9
80	32.8	67.2	1.605	-	0.45	3.8

with an RMS surface roughness of 0.54 nm had the lowest loss of 0.1 dB/cm at 638 nm. The loss in these films appears to decrease at higher wavelengths with all of them measuring a loss at 0.1 dB/cm at a wavelength of 980 nm. Given that scattering is scaled as  $\lambda^{-4}$ , we would expect the loss decrease at higher wavelengths if the primary loss mechanism was caused through light scattering at the interfaces.

The RMS roughnesses of the as-deposited films are less than those of similar  $\text{GeO}_2$  films grown by RF magnetron sputtering [23] and DC reactive sputtering [22]. Thin films grown using PARMs have been shown to exhibit high packing densities and film surface uniformity [38]. The losses in these  $\text{GeO}_2$  films are also shown to be less than the films grown using other methods without the need for post-processing.

#### 4. Conclusions

We have demonstrated low-loss  $\text{GeO}_2$  thin films fabricated by PARMs AC sputtering and show an improvement in losses for the as-deposited  $\text{GeO}_2$  films compared to previous deposition techniques. The films were shown to be stoichiometric through RBS measurements, and the refractive index was modeled through VASE and prism coupler measurements and determined to be around 1.605 at 638 nm. The film composition and refractive index were found to be constant for different O<sub>2</sub> flow rates. Small amounts of hydrogen were present in the films, which might not be considered detrimental as hydrogen is known to sensitize  $\text{GeO}_2$  to change its refractive index through the photorefractive effect. The films were shown to have sub-nanometer surface roughnesses down to 0.5 nm, as determined through AFM, and losses as low as 0.1 dB/cm at 638 nm without the need for post-deposition process steps. The films also had high uniformity across a 3-inch wafer with a relative thickness deviation of 0.8%. The low surface roughnesses and losses demonstrate ion-assisted dual AC sputtering to be especially attractive for the deposition of optical waveguide materials.

#### CRediT authorship contribution statement

J.W. Miller: Methodology, Formal analysis, Writing - original draft, Writing - review & editing. M. Chesaux: Methodology, Writing - original draft. D. Deligiannis: Resources, Funding acquisition. P. Mascher: Supervision, Writing - review & editing. J.D.B. Bradley: Supervision, Writing - review & editing.

#### Declaration of Competing Interests

None.

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